

Short-range magnetic order and temperature-dependent properties of cupric oxide

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2010 J. Phys.: Condens. Matter 22 045502

(<http://iopscience.iop.org/0953-8984/22/4/045502>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 129.252.86.83

The article was downloaded on 30/05/2010 at 06:38

Please note that [terms and conditions apply](#).

Short-range magnetic order and temperature-dependent properties of cupric oxide

X Rocquefelte¹, M-H Whangbo², A Villesuzanne³, S Jobic¹,
F Tran⁴, K Schwarz⁴ and P Blaha⁴

¹ Institut des Matériaux Jean Rouxel, Université de Nantes, CNRS, UMR6502,
Boîte Postale 32229, 44322 NANTES Cedex 3, France

² Department of Chemistry, North Carolina State University, Raleigh, NC 27695-8204, USA

³ Institut de Chimie de la Matière Condensée de Bordeaux, UPR9048, ICMCB-CNRS,
av. Dr A Schweitzer, 33608 Pessac Cedex, France

⁴ Institute of Materials Chemistry, Vienna University of Technology, Getreidemarkt 9/165-TC,
A-1060 Vienna, Austria

Received 24 September 2009, in final form 9 November 2009

Published 5 January 2010

Online at stacks.iop.org/JPhysCM/22/045502

Abstract

The temperature dependence of the optical and magnetic properties of CuO were examined by means of hybrid density functional theory calculations. Our work shows that the spin exchange interactions in CuO are neither fully one-dimensional nor fully three-dimensional. The temperature dependence of the optical band gap and the ⁶³Cu nuclear quadrupole resonance frequency of CuO originate from the combined effect of a strong coupling between the spin order and the electronic structure and the progressive appearance of short-range order with temperature.

(Some figures in this article are in colour only in the electronic version)

In strongly correlated transition-metal oxides the coupling between spin, charge, orbital and vibrational degrees of freedom favours the emergence of exotic properties such as colossal magnetoresistance, high-temperature superconductivity or multiferroicity. Cupric oxide CuO has been extensively studied since the discovery of cuprate superconductors due to its close resemblance in structure and magnetic properties. In particular, some features of cuprate superconductors appear also in CuO, for example, Zhang–Rice singlets at the top of its valence band [1], charge stripes [2] and a strong magneto–lattice coupling [3]. Recently, CuO has been found to be multiferroic at high temperature [4].

The optical band gap and the absorption edge of CuO exhibit a significant dependence on temperature as shown by Marabelli *et al* [5]. They concluded that the electronic states involved in the optical absorption processes have a localized character, originating from a strong electron–lattice coupling. Above $T_N = 230$ K, the three-dimensional (3D) magnetic ordering temperature, CuO is usually described as a quasi-one-dimensional (1D) Heisenberg antiferromagnet with a strong spin exchange (i.e. $J = -67 \pm 20$ meV) [6–9]. Thus, one

may wonder whether the strong temperature dependence of the optical band gap and the absorption edge in CuO arise from a strong coupling of the electronic structure with the spin order rather than with the lattice. In this paper we show that this is indeed the case by examining the spin exchange interactions, the dielectric function, the band gap and the electric field gradients (EFGs) of CuO on the basis of first principles density functional theory (DFT) [10] electronic structure calculations.

The structure of CuO can be viewed as built upon corner- and edge-sharing square-planar CuO₄ units, which form $(-\text{Cu}-\text{O}-)_{\infty}$ zigzag chains running along the [101] and $[10\bar{1}]$ directions of the chemical unit cell (a, b, c), i.e. the a' and c' directions of the magnetic cell, respectively (see figure 1). The Cu²⁺ spin moments order to have an incommensurate antiferromagnetic (AFM) structure with propagation vector $q = (0.506, 0, -0.483)$ below $T_N = 230$ K, and subsequently a commensurate structure with $q = (0.5, 0, -0.5)$ below $T_L = 213$ K [6–8], in which Cu²⁺ spins have AFM and ferromagnetic (FM) arrangements along the $[10\bar{1}]$ and [101] directions, respectively [7]. The dimensionality of the magnetic structure has been controversial; the magnetic,

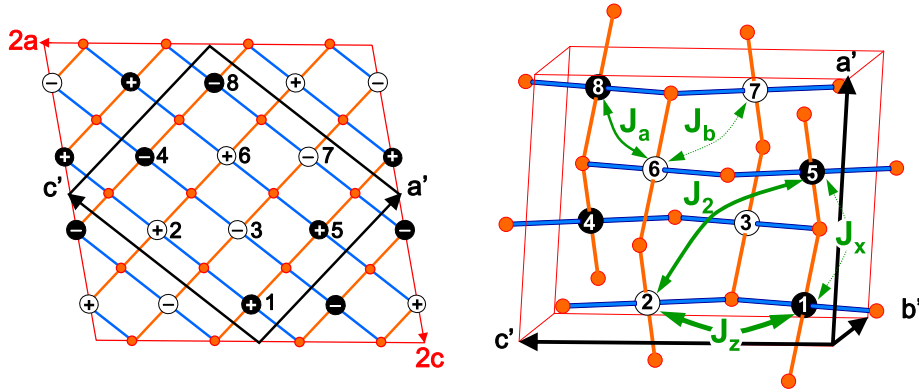


Figure 1. Schematic representations of the magnetic structure of CuO observed below $T_L = 213$ K. The left panel shows a projection along the b axis of the $2a \times 2c$ supercell containing 16 Cu^{2+} ions. The magnetic unit cell (a' , b' , c'), shown in black, contains eight Cu^{2+} ions. The plus and minus signs indicate the position of copper atoms along the b' axis. The right panel shows a 3D perspective view of the magnetic structure. The O atoms are presented by orange (grey) circles and the Cu^{2+} sites along the b -direction are depicted as filled and empty circles representing up-spin and down-spin, respectively. The definitions of the five spin exchange paths J_z , J_2 , J_a , J_b and J_x in the magnetic unit cell (a' , b' , c') are also given.

thermodynamic and neutron data of CuO above T_N can be reproduced using either a 1D AFM chain or a two-dimensional (2D) AFM lattice model. This picture of low-dimensional magnetism is supported by spin dimer analysis based on extended Hückel tight binding (EHTB) calculations [11, 12], according to which CuO can be described as having 1D AFM chains. However, in a recent pseudopotential self-interaction-free (pseudo-SIC) DFT study, Filippetti and Fiorentini [13] reported that CuO is an antiferromagnet with fully 3D spin exchange interactions, and the magnetic orbital of each square-planar CuO_4 unit is not the Cu $d_{x^2-y^2}$ but the Cu d_{z^2} orbital. The latter conclusions are inconsistent with the observed low-dimensional magnetic properties of CuO and the well-established fact that the highest-lying d-block orbital of a square-planar complex of a transition-metal element is not represented by the d_{z^2} but by the $d_{x^2-y^2}$ orbital [12, 14]. Therefore, an accurate and independent determination of the spin exchange interactions is highly desirable.

To evaluate the spin exchange interaction of CuO, we initially considered the nine ordered spin states (i.e. AF_x , AF_y , AF_z , AF_a , AF_{-101} , FM, FM_1 , FM_x and FM_z) depicted in figure 2 assuming collinear arrangements of magnetic moments. These ordered spin states can be constructed by using the magnetic unit cell consisting of eight formula units (FUs) shown in figure 1.

DFT calculations were performed for these states on the basis of the room-temperature (RT) single-crystal x-ray structure [15]⁵ using the WIEN2k program package [17]. The present calculations employed a hybrid functional based on the Perdew–Burke–Ernzerhof (PBE) [18] functional of the generalized gradient approximation (GGA). In hybrid functionals [19], the exchange energy is obtained by mixing Hartree–Fock (HF) and DFT exchange energies in a certain proportion. The optimal fraction α of the Hartree–Fock exchange varies according to the systems under examination (nature of the elements, local arrangement,

and dimensionality), their properties and the exchange–correlation functionals employed in the mixing [20–22]. The implementation of this approach in solid state DFT codes is quite recent, but the relevance in describing transition-metal oxide [22–24] and excited state properties (e.g. optical band gap, spin exchange, etc) has been demonstrated in several studies [23–26]. In particular, the solutions provided by such hybrid calculations go beyond the Kohn–Sham ground-state solutions, allowing time-consuming many-body calculations to be quite accurately reproduced [25]. In addition, it should be noticed that the present results for excited state properties (magnetism, optics) are corroborated by ground-state property results like the electric field gradient (EFG), which was shown to be extremely sensitive to the treatment of the correlated electrons, as shown in [27]. To ensure the validity of the present hybrid generalized gradient approximation (hybrid/GGA) calculations, and more specifically the value of the mixing parameter, $\alpha = 0.15$, we have considered the low-temperature AFM arrangement of CuO. Particularly, ground-state properties like the value of the EFG of the Cu atom of $-8.50 \times 10^{21} \text{ V m}^{-2}$ (see below) and the spin moments on the Cu and O atoms of 0.65 and 0.12 μ_B , respectively, are in good agreement with the experimental values of $-7.80 \times 10^{21} \text{ V m}^{-2}$, 0.65 and 0.14 μ_B [7], respectively. In addition, our calculations lead to an optical band gap of about 1.42 eV, which compares well with the experimental value of 1.4 eV [28]. While our results are in reasonable agreement with previous DFT calculations, their values for the band gap and Cu spin moments of 1.9 eV and 0.74 μ_B in LDA + U ($U_{\text{eff}} = 6.5$ eV) [29] and of 2.2 eV and 0.72 μ_B in pseudo-SIC approaches [13] were significantly overestimated.

The relative energies of the nine ordered spin states obtained from our hybrid/GGA calculations are summarized in table 1. To extract the spin exchange parameters J_a , J_b , J_x , J_z and J_2 of CuO as defined in figures 1 and 2, we express the total spin exchange interactions of the nine ordered spin states in terms of the spin Hamiltonian $\hat{H} = -\sum_{i < j} J_{ij} \hat{S}_i \cdot \hat{S}_j$, where J_{ij} ($= J_a, J_b, J_x, J_z$ or J_2) refers to the spin exchange

⁵ Our calculations for the single-crystal structure of CuO determined at 11 K [16] provide almost identical results.

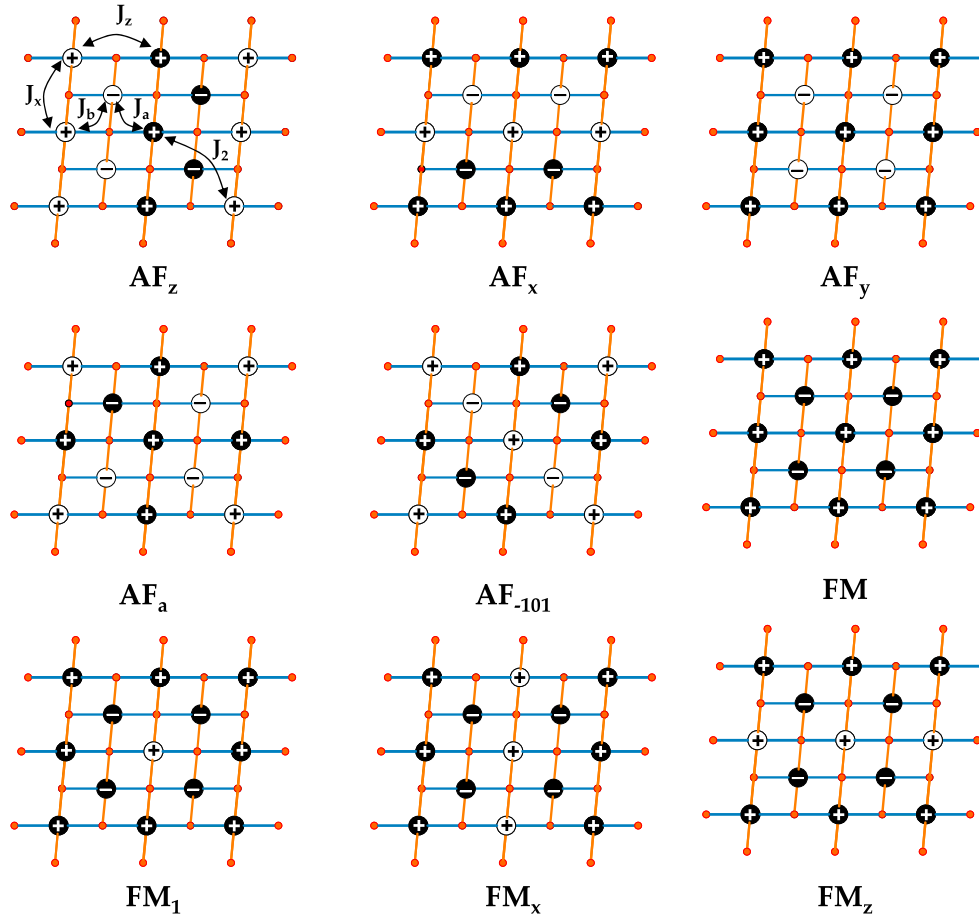


Figure 2. Schematic projection views, along the b' axis, of the nine ordered spin arrangements of the Cu^{2+} ions in CuO .

parameter for the spin sites i and j .⁶ With this convention, an AFM spin exchange interaction is represented by a negative J value. Thus by mapping the energy differences between the ordered spin states given by the spin Hamiltonian onto those given by hybrid/GGA calculations, we obtain the values of the five spin exchange parameters listed in table 2.

In agreement with the neutron diffraction study [7] our calculations show that the spin arrangement AF_z is the most stable magnetic state. In the study of Filippetti and Fiorentini [13] all spin exchange interactions are comparable in magnitude (i.e. $J_z \approx -2J_x \approx -3J_b \approx 3J_2 \approx -5J_a$), and only J_z and J_2 favour AFM interactions. However, the present hybrid/GGA calculations (with $\alpha = 0.15$) show that only three interactions are significant ($J_z \approx 4J_2 \approx 7J_a$) and favour AFM interactions. J_z is the strongest spin exchange interaction, as already found in the EHTB study ($J_z \approx 16J_x \approx 19J_b \approx 19J_2 \approx 310J_a$) [11]. Table 1 shows that only three terms (J_z , J_2 and J_x) contribute to the stability of the AF_z state. J_x was previously considered as the driving force of the FM couplings along x [13], but is very weak. Thus, the present calculations evidence that the FM layers of the AF_z state are a direct consequence of the fact that the 1D AFM chains (Cu1–Cu2 and

⁶ There are three different Cu–O...O–Cu super-superexchange paths contributing to J_2 , namely, J_{2a} , J_{2b} and J_{2c} with Cu...Cu distances of 5.800, 5.129 and 4.680 Å, respectively. Our additional calculations show that J_{2a} and J_{2c} are negligibly small.

Table 1. Energy expressions per Cu atom and the corresponding relative energies (Note: the reference state, with respect to which the relative energies are given, is defined as the hypothetical state for which the total spin exchange energy is zero. For example, $E(\text{reference}) \equiv [2E(\text{FM}_1) - E(\text{FM})]/2$) (in meV) deduced from hybrid/GGA calculations (α is the fraction of Hartree–Fock exchange).

State	Energy expression	$\alpha = 0.15$	$\alpha = 0.25$
AF_z	$(-J_x + J_z + J_2)/4$	-40.4	-25.8
AF_{-101}	$(-J_a + 2J_b + J_x + J_z - J_2)/4$	-17.4	-11.7
AF_a	$+J_a/8$	-2.3	-0.5
FM_x	$-J_x/4$	-0.3	-0.7
FM_1	$(-J_a - 2J_b - J_x - J_z - J_2)/8$	20.7	11.7
AF_x	$(+J_x - J_z + J_2)/4$	25.3	16.0
FM_z	$-J_z/4$	32.0	20.1
AF_y	$(+J_a + 2J_b - J_x - J_z - J_2)/4$	36.6	25.0
FM	$(-J_a - 2J_b - J_x - J_z - J_2)/4$	41.5	23.5

Cu5–Cu6 in figure 1) made up of the dominant superexchange interaction J_z are also antiferromagnetically coupled through the Cu–O...O–Cu super-superexchange interaction J_2 (Cu2–Cu5 and Cu1–Cu6).

Given the experimental estimate of $J = -67 \pm 20$ meV [6–9], the exchange parameters obtained from the hybrid/GGA calculations with $\alpha = 0.15$ are overestimated (i.e. $J_z = -128.8$ meV, table 2). These small interaction

Table 2. Geometrical and spin exchange J (in meV) parameters of CuO obtained from the hybrid/GGA calculations with $\alpha = 0.15$ and $\alpha = 0.25$.

	J_a	J_b	J_x	J_z	J_2 (see footnote 6)
$d_{\text{Cu-Cu}}$ (Å) [15]	2.901	3.083	3.173	3.749	5.129
$\angle\text{Cu-O-Cu}$ (deg) [15]	95.71	104.02	108.91	145.81	
Calculated J with $\alpha = 0.15$	-18.2	+4.2	+2.6	-128.8	-30.1
Calculated J with $\alpha = 0.25$	-4.0	+3.5	+3.0	-80.5	-19.6

energies are very sensitive to the chosen DFT functional. As expected, if we use a larger mixing value of $\alpha = 0.25$ (PBE0 hybrid functional) [20, 21], we obtain smaller J values in much better agreement with the experimental estimation. In particular, the strongest interaction $J_z = -80.5$ meV, is then in good agreement with the reported experimental value. However, the larger fraction of HF exchange leads to a significant overestimation of the optical band gap of about 2.4 eV, the EFG of about -11.96×10^{21} V m $^{-2}$ and the Cu spin moment of $0.74 \mu_B$, while the O spin moment is then underestimated, $0.09 \mu_B$. As a matter of fact, with a larger mixing value the Cu-d states are more localized, leading to a reduced orbital overlap with O-p states. Consequently the ionicity of the Cu–O chemical bond is increased, leading to a relocalization of the magnetic moment on the Cu site, an increase of the band gap, and a reduction of the spin exchange mediated by the orbital overlap between Cu and O atoms.

We should keep in mind that the choice of the mixing parameter in such mean-field approaches is a matter of compromise: higher values for spin exchange estimation, but smaller values for spectroscopic data simulation. In any case, the hybrid/GGA calculations with both $\alpha = 0.15$ and 0.25 give the same trends in the spin exchange parameters, namely, $J_z/J_2 \approx 4$ and other J values are much smaller. These findings are in strong disagreement with previous calculations by Filippetti and Fiorentini [13] and restore the previously accepted view of magnetic interactions in CuO [6–9, 11]. The above picture of CuO, deduced from our hybrid/GGA calculations, predicts that the spin correlation lengths along the c' - and a' -directions slightly above T_N should be controlled by J_z and J_2 , respectively. Consistent with this prediction, the observed spin correlation length along the c' -direction is about 3.5 times greater than that along the a' -direction (i.e. 700 versus 200 Å) [30], which is consistent with the fact that $J_z/J_2 \approx 4$.

From the viewpoint of the calculated spin exchange parameters, the temperature dependence of the magnetic susceptibility of CuO, which shows a broad maximum at around $T_{\text{max}} = 550$ K [6–9], can be described as follows: at high temperature (around T_{max} and above), CuO exhibits short-range 1D AFM order based on J_z . As the temperature is lowered, CuO undergoes short-range 2D AFM order based on J_z and J_2 (above T_N and below T_{max}), which is followed by incommensurate 3D AFM long-range order based largely on J_z , J_2 and J_a (below T_N and above T_L) and then by commensurate 3D AFM long-range order (below T_L). As can be seen from figure 1 and table 1, both J_a and J_b do not contribute to the stability of the AF $_z$ phase but can introduce spin frustration between adjacent 2D rectangular nets

defined by J_z and J_2 . The latter might be responsible for the incommensurate 3D AFM long-range order between 213 and 230 K, in which every second 1D chain made up of J_z undergoes a spiral spin order responsible for the ferroelectric polarization [4]. The occurrence of short-range magnetic order is consistent with the observations that no paramagnetic scattering is detected down to 550 K in the paramagnetic regime ($T > T_N$) [7], and that more than 70% of the spins are ordered above T_N according to the specific-heat measurements [31].

Let us now explore the temperature dependence of the optical properties of CuO. By considering low- and high-temperature crystal structures of CuO (see footnote 5), we note that the small changes in the crystal structure with temperature have a negligible effect on the calculated dielectric function ϵ , and then on the temperature dependence of the optical band gap of CuO [5]. Thus, the band gap decrease when temperature is raised is not related to a lattice expansion effect. In normal semiconductors (e.g. Si) the decrease of the band gap with temperature can be well described by thermal expansion and electron–phonon interactions [32]. In Si the gap variation from 26 to 305 K is only about 0.037 eV, whereas in CuO this variation is an order of magnitude larger and thus a mechanism other than electron–phonon interaction needs to be invoked. In contrast, the temperature-induced rearrangement of the spin order, described above, may influence the electronic structure and its related optical transitions. Upon raising the temperature above T_N , the thermal energy induces spin excitations, which destroy the 3D long-range AFM order, but short-range 2D/1D AFM order remains. However, it is important to note that above the Néel temperature, not only the spin dynamics is changed and the spin correlation is destroyed, but also the spin symmetry is different (AFM spin symmetry is lost). In the case of low-dimensional magnetic systems with one strong J value, such as in CuO, the loss of the anti-parallel correlation between the spins allows them to align along the magnetic field. This change in symmetry is progressive and explains the continuous increase in the magnetic susceptibility from T_N to $T_{\text{max}} = 550$ K, when a magnetic field is applied in magnetic susceptibility measurements. Above T_{max} the spin correlation is expected to be negligible and then the spins are in a paramagnetic regime.

To model the high-temperature behaviour of the magnetic properties of CuO, one should progressively destroy the AFM spin correlation, i.e. by introducing a paramagnetic (PM) component in CuO. Such treatment requires the use of dynamical mean-field theory approaches [33], at least. However, in the present case we are interested in simulating properties which are very sensitive to changes of the local

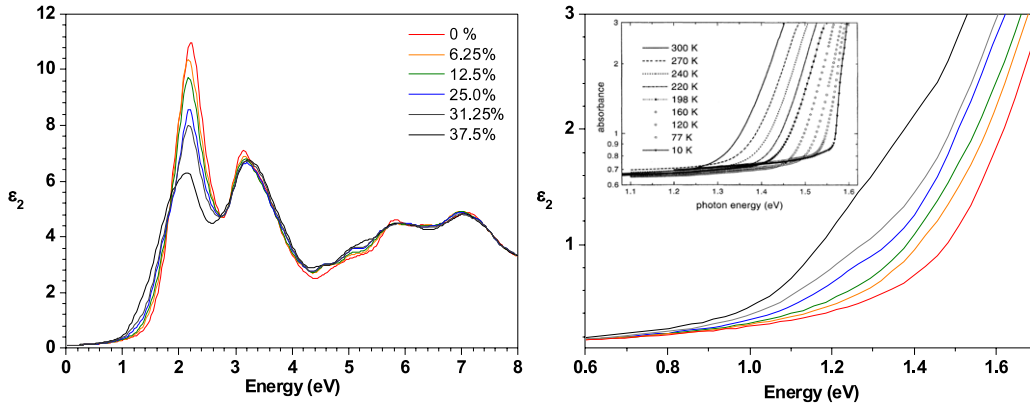


Figure 3. Isotropic imaginary part ε_2 of the dielectric function of CuO as a function of the incident light energy. The left panel shows the ε_2 curves obtained from the hybrid/GGA calculations with $\alpha = 0.15$ for the magnetic states of CuO with $P = 0, 6.25, 12.5, 25.0, 31.25$ and 37.5% . The right panel shows a zoomed-in view around the absorption edge, where the inset shows the experimental absorbance of CuO at the absorption edge as a function of temperature [5].

environment of the probed atom. This can also be simulated by a static approximation. The local symmetry of the spin density has a strong impact on the orbital overlap and then on the related electronic structure. For these reasons, we have chosen to focus our attention on the temperature effects on the local symmetry of the spin density and its related consequences on the optical properties and EFG, neglecting the dynamical aspect (i.e. the spin dynamics). To properly describe the evolution of the magnetic structure of CuO above T_N , one should describe the emergence of short-range spin ordered states by progressively introducing PM domains in the AFM state. However, it is interesting to note that FM and PM states have the same symmetry, in contrast to the AFM state. Our strategy was therefore to mimic the temperature effect on the spin density symmetry by flipping statically the spin on certain atoms, introducing local FM domains in the AFM state. Although these static models do not reproduce the reality of the magnetic structure of CuO at high temperature, they allow us to properly describe the symmetry changes that are crucial to describe electronic structure related properties.

The electronic structures for various short-range spin ordered states of CuO, which occur when the temperature is raised, were simulated by hybrid/GGA calculations for CuO with a $(4a \times b \times 4c)$ supercell containing 64 FUs. The five short-range spin ordered models we have envisioned are all constructed based on the AF_z state, for which the 64 FUs supercell contains 64 AFM interactions along the c' -direction. For convenience, these five short-range spin ordered states of CuO will be labelled in terms of the percentage P of destroyed AFM interactions within the AFM chains of Cu^{2+} ions along the c' -direction. For instance, $P = 0\%$ describes the AF_z state, while $P = 6.25\%$ describes short-range ordered states in which 6.25% of AFM interactions along the c' -direction have been destroyed. All of the five models, i.e. for $P = 6.25, 12.5, 25, 31.25, 37.5\%$, have been constructed in such a way that the total magnetic moment in the unit cell is equal to zero (to reproduce the situation corresponding to the PM state). Figure 3 compares the imaginary parts of the dielectric function ε_2 (related to the optical absorption)

calculated for the states with $P = 0, 6.25, 12.5, 25, 31.25, 37.5\%$. Note that in our model P is expected to increase with increasing temperature, but of course never reaches an FM state (which is different to a $P = 100\%$ state, which exhibits a zero total magnetic moment). The predicted optical response corresponding to a hypothetical FM state is also given in figure 4.

As can be seen in figures 3 and 4, the RT experimental ε_2 curve [34, 35] shows two broad absorption peaks. The latter are not reproduced by the electronic structure of either the AF_z or the FM state. However, the ε_2 curves calculated for the electronic states with short-range spin order show an optical response intermediate between those of the AF_z and FM states, hence reproducing the essential feature of the experimental ε_2 curve. Note that the lower-energy absorption is overestimated by the AF_z state while it is absent in the FM state. Thus, figure 3 reveals that, with increasing temperature (i.e. with decreasing the number of AFM interactions within each chain of J_z), the lower-energy absorption peak undergoes three significant changes; a reduction in the intensity, a displacement toward the lower energy and a decrease in the steepness of the absorption edge. The latter two findings explain the large temperature dependence (from 10 to 300 K) of the experimentally observed absorption edge [5], namely, a gradual decrease in the band gap by about 0.25 eV and in the steepness of the absorption edge. Our simulations show that the optical band gap is reduced by about 0.3 eV when P increases from 0% to 37.5%. The present study clearly evidences that the temperature dependence of the optical properties of CuO is a consequence of strong coupling between the spin order and the electronic structure.

To better understand the origin of the strong coupling between the optical response and the degree of spin order in CuO, we compare the density of states (DOS) calculated for the $P = 0, 37.5\%$ and hypothetical FM states (figure 4). In all cases the absorption edge is mainly due to a charge transfer transition from the O 2p to the empty Cu $3d_{x^2-y^2}$ bands. Both the majority and the minority spins are involved in this transition for the $P = 0$ and 37.5% states, but only

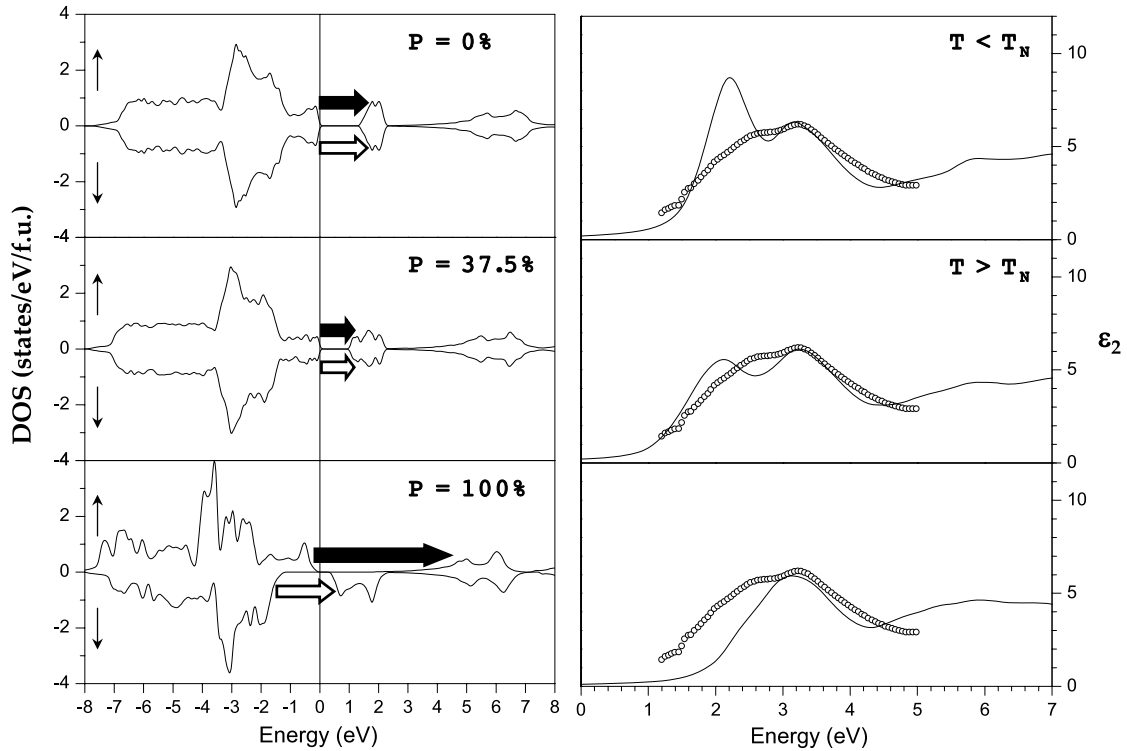


Figure 4. DOS plots (left panel) and ϵ_2 curves (right panel) obtained from the hybrid/GGA calculations with $\alpha = 0.15$ for the $P = 0, 37.5$ and 100% states of CuO. In the right panel, the experimental ϵ_2 plot obtained at room temperature is given as empty circles [34].

the minority spins contribute for the hypothetical FM state. The latter explains why the low-energy absorption is absent in the FM state. The DOS plots show that the O 2p and Cu 3d states are strongly affected on going from the AF_z state to the $P = 37.5\%$ state and then to the FM state. The 3d_{x²-y²} bands progressively increase in width and hence induce a reduction in the optical gap when the temperature is raised (figure 4). It should be emphasized that although the experimental ϵ_2 data appear featureless, they were obtained in two independent studies [34, 35]. The broad feature of ϵ_2 at RT is not a consequence of experimental limitation, but reflects a broadened DOS at RT that is brought about by the progressive loss of AFM interactions within each chain of J_z .

The essence of the above observations can be readily accounted for by considering the electronic structure of a spin-1/2 uniform AFM chain made up of the spin exchange J_z . If this chain is described by the nearest-neighbour hopping integral t and the on-site repulsion U , the up- and down-spin bands have a width proportional to t in the FM state, but t^2/U in the AFM state. For a strongly AFM chain, $t \gg t^2/U$, so that the bandwidth is significantly narrower in the AFM state than in the FM state, as has been pointed out by Xiang *et al* in their analysis of the dependence of the electronic structure of SrFeO₂ on its spin arrangement [36]. Here, we should not forget that FM segments (two or three atoms with parallel spin alignment) have been introduced in the system to mimic the local symmetry changes when PM segments are created due to the temperature increase. For symmetry reasons, the bandwidths of PM and FM states will be larger than the bandwidth of the AFM state. Because of the similarity of

PM and FM states in terms of spin density symmetry, our calculations correctly reproduce the change of spin-orbital overlap due to the incorporation of PM domains in the AFM chain of J_z . From the present calculations, it could then be deduced that the temperature increase leads to a progressive replacement of the AFM interactions in the chain of J_z by PM domains, which generates discrete energy levels, the high- and the low-lying levels of which would lie above and below the narrow band associated with the AFM regions of the chain, respectively. This explains why the loss of 3D AFM long-range order in CuO decreases the band gap and the steepness of the absorption edge and why the states participating in the excitation have a localized character.

Finally, we show that the temperature dependence of the ⁶³Cu nuclear quadrupole resonance (NQR) frequency in CuO is another signature of strong coupling between the spin order and the electronic structure in CuO. The ⁶³Cu NQR frequency of CuO increases linearly with temperature [37, 38], and this finding is not explained by considering the lattice expansion with increasing temperature. The ⁶³Cu NQR frequency is proportional to the EFG and is very sensitive to small changes in the electronic structure. The strong coupling between the spin order and the electronic structure discussed above should affect the electron density around the Cu atoms in CuO, and hence the electric field gradient (EFG) of the Cu atoms. We calculated the EFGs of the Cu atoms by performing hybrid/GGA calculations with $\alpha = 0.15$ for three states, namely the low-temperature 3D long-range AF_z spin order ($P = 0\%$), the hypothetical 3D long-range FM spin order ($P = 100\%$) and the 2D/1D short-range spin order state, $P = 37.5\%$,

which leads to the best agreement with the RT experimental optical data. Our calculations show that the EFGs of the Cu atoms increase almost linearly with increasing P . For example, the calculated EFG is $-8.50 \times 10^{21} \text{ V m}^{-2}$ for the AF_z state ($P = 0\%$), $-8.78 \times 10^{21} \text{ V m}^{-2}$ for the $P = 37.5\%$ state and $-9.28 \times 10^{21} \text{ V m}^{-2}$ for the FM state. Thus, our calculations predict that, on going from the low-temperature ($P = 0\%$) to the high-temperature ($P = 37.5\%$) magnetic structure, the EFG should increase by about 3%. This is consistent with the experimental observation between 77 and 295 K (EFG = -7.80×10^{21} and $-8.04 \times 10^{21} \text{ V m}^{-2}$, respectively, by using $Q = -0.211$ barn for ^{63}Cu).

In summary, the temperature-induced change in the dimensionality of the magnetic structure of CuO from 1D to 3D is well explained in terms of the spin exchange parameters deduced from the present work. The temperature dependence of the magnetic susceptibility, optical band gap and ^{63}Cu NQR frequency in CuO originates from a strong coupling between the spin order and the electronic structure. It remains for us to explore whether the present theoretical scenario could be generalized to other highly correlated transition-metal oxides with strong J values, such as the high- T_c cuprate parent compounds.

Acknowledgments

The authors thank the Centre de Calcul Intensif des Pays de la Loire for generous computing resources. The work at NCSU was supported by the Office of Basic Energy Sciences, Division of Materials Sciences, US Department of Energy, under Grant DE-FG02-86ER45259. PB and FT were supported by the Austrian Science Fund P20271-N17. XR and SJ would like to thank C Payen for helpful discussions.

References

- [1] Tjeng L H, Sinkovic B, Brookes N B, Goedkoop J B, Hesper R, Pellegrin E, de Groot F M F, Altieri S, Hulbert S L, Shekel E and Sawatzky G A 1997 *Phys. Rev. Lett.* **78** 1126
- [2] Zheng X G, Xu C N, Tomokiyo Y, Tanaka E, Yamada H and Soejima Y 2000 *Phys. Rev. Lett.* **85** 5170
- [3] Yamada H, Zheng X-G, Soejima Y and Kawaminami M 2004 *Phys. Rev. B* **69** 104104
- [4] Kimura T, Sekio Y, Nakamura H, Siegrist T and Ramirez A P 2008 *Nat. Mat.* **7** 291
- [5] Marabelli F, Parravicini G B and Salghetti-Drioli F 1995 *Phys. Rev. B* **52** 1433
- [6] O'Keefe M and Stone F S 1962 *J. Phys. Chem. Solids* **23** 261
- [7] Forsyth J B, Brown P J and Wanklyn B M 1988 *J. Phys. C: Solid State Phys.* **21** 2917
- [8] Yang B X, Thurston T R, Tranquada J M and Shirane G 1989 *Phys. Rev. B* **39** 4343
- [9] Shimizu T, Matsumoto T, Goto A, Chandrasekhar Rao T V, Yoshimura K and Kosuge K 2003 *Phys. Rev. B* **68** 224433
- [10] Kohn W and Sham L J 1965 *Phys. Rev.* **140** A1133
- [11] Dai D, Koo H-J and Whangbo M-H 2004 *Inorg. Chem.* **43** 4026
- [12] Whangbo M-H, Koo H-J and Dai D 2003 *J. Solid State Chem.* **176** 417
- [13] Filippetti A and Fiorentini V 2005 *Phys. Rev. Lett.* **95** 086405
- [14] Albright T A, Burdett J K and Whangbo M-H 1985 *Orbital Interactions in Chemistry* (New York: Wiley) chapter 16
- [15] Asbrink S and Norrby L-J 1970 *Acta Crystallogr. B* **26** 8
- [16] Brese N E, O'Keefe M, Ramakrishna B L and Von Dreele R B 1990 *J. Solid State Chem.* **89** 184
- [17] Blaha P, Schwarz K, Madsen G K H, Kvasnicka D and Luitz J 2001 *WIEN2k: An Augmented Plane Wave + LO Program for Calculating Crystal Properties* (Vienna: TUW)
- [18] Perdew J P, Burke K and Ernzerhof M 1996 *Phys. Rev. Lett.* **77** 3865
- [19] Becke A D 1993 *J. Chem. Phys.* **98** 5648
- [20] Ernzerhof M and Scuseria G E 1999 *J. Chem. Phys.* **110** 5029
- [21] Adamo C and Barone V 1999 *J. Chem. Phys.* **110** 6158
- [22] Cora F, Alfredsson M, Mallia G, Middlemiss D, Mackrodt W C, Dovesi R and Orlando R 2004 *Struct. Bond.* **113** 171
- [23] Tran F, Blaha P, Schwarz K and Novák P 2006 *Phys. Rev. B* **74** 155108
- [24] Francini C, Podloucky R, Paier J, Marsman M and Kresse G 2007 *Phys. Rev. B* **75** 195128
- [25] Paier J, Asahi R, Nagoya A and Kresse G 2009 *Phys. Rev. B* **79** 115126
- [26] Moreira I P, Calzado C J, Malrieu J-P and Illas F 2006 *Phys. Rev. Lett.* **97** 087003
- [27] Blaha P, Schwarz K and Novák P 2005 *Int. J. Quantum Chem.* **101** 550
- [28] Ghijsen J, Tjeng L H, van Elp J, Eskes H, Westerink J, Sawatzky G A and Czyzyk M T 1988 *Phys. Rev. B* **38** 11322
- [29] Anisimov V I, Zaanen J and Andersen O K 1991 *Phys. Rev. B* **44** 943
- [30] Chattopadhyay T, McIntyre G J, Vettier C, Brown P J and Forsyth J B 1992 *Physica B* **180/181** 420
- [31] Loram J W, Mirza K A, Joyce C P and Osborne A J 1989 *Europhys. Lett.* **8** 263
- [32] Lautenschlager P, Garriga M, Vina L and Cardona M 1997 *Phys. Rev. B* **36** 4821
- [33] Held K 2007 *Adv. Phys.* **56** 829
- [34] Ito T, Yamaguchi H, Masumi T and Adachi S 1998 *J. Phys. Soc. Japan* **67** 3304
- [35] Babonas G J, Reza A, Maltsev V, Galickas A and Dapkus L 2002 *Lith. J. Phys.* **42** 161
- [36] Xiang H J, Wei S-H and Whangbo M-H 2008 *Phys. Rev. Lett.* **100** 167207
- [37] Itoh Y, Imai T, Shimizu T, Tsuda T, Yasuoka H and Ueda Y 1990 *J. Phys. Soc. Japan* **59** 1143
- [38] Graham R G, Fowler D, Lord J S, Riedi P C and Wanklyn B M 1991 *Phys. Rev. B* **44** 7091